

Family Resemblances among QCD's Close Relatives

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Outline

- Large N QCD – why it's interesting to a lattice person
- Mesonic observables
- Baryons and their rotor spectrum
- Flavor $SU(3)$ spectroscopy
- Alternate large N constructions

See Phys. Rev. D86, 034508 (arXiv:1205.0235); D89, 014506
(arXiv:1308.4114; D90, 014505 (arXiv:1404.2301)

Just a fun project for the “post-QCD” era

Why is large- N lattice QCD interesting?

- It's the source of all qualitative understanding of hadronic physics
- “No free parameters for QCD” since the coupling runs.
- 't Hooft (1974): Treat $1/N$, where N is the number of colors, as a small number
- Familiar scaling story for mesons as $\bar{q}q$ objects with masses independent of N
- Extensive successful (pre-lattice) phenomenology for matrix elements (B_K , etc) and their regularities
- An interesting set of stories for baryons, too
- Since 1974, MANY other large- N stories
 - Higher representation fermions – Corrigan - Ramond, AS2 rep fermions
 - Orientifold equivalence, volume reduction, etc etc
- Composite Higgs phenomenology often uses large- N to extrapolate predictions

And now we can simulate these systems, rather easily, and see how well large- N worked.

Fun, easy projects!

Baryons in large N

- Mass $\propto N$ (for fundamentals, to $N(N - 1)/2$ for AS2)
- Hartree approximation: 2-body interaction $V \sim g^2/N$, many little interactions, central limit theorem
 - baryon size is N independent
- Large- N QCD is EFT of mesons
 - Coupling is $1/F_\pi^2 \sim 1/N$ (fundamentals)
 - Baryon mass $\sim N$ or $1/\text{coupling}$ – soliton-like – Skyrme model as an explicit phenomenology
- Either way, mass formulas in $1/N$ – tested
- Spin - isospin locking for flavor $SU(2)$: $J = I$ for $J = N/2, N/2 - 1, \dots$ – tested
- Mass relations imply stringent matrix element relations and vice versa – so far, not tested

Technical issues for lattice simulations

I am using an NCOL version of the Milc code written by Svetitsky, Shamir and me for BSM studies

- $N = 3, 5, 7$ with fundamentals
- I did quenched approximation, first
 - Easy to justify for large N
 - Have since done dynamical fermion sims for $N = 4$ AS2
- Tune bare couplings to match lattice spacings (from the potential)
- Observe this matches $\lambda = g^2 N$, too
- Match to same volumes, roughly same quark masses
- Compare the usual lattice-y dimensionless ratios of things
- Need baryon correlators for arbitrary N – a long story with much room for improvement
- Other little technical issues

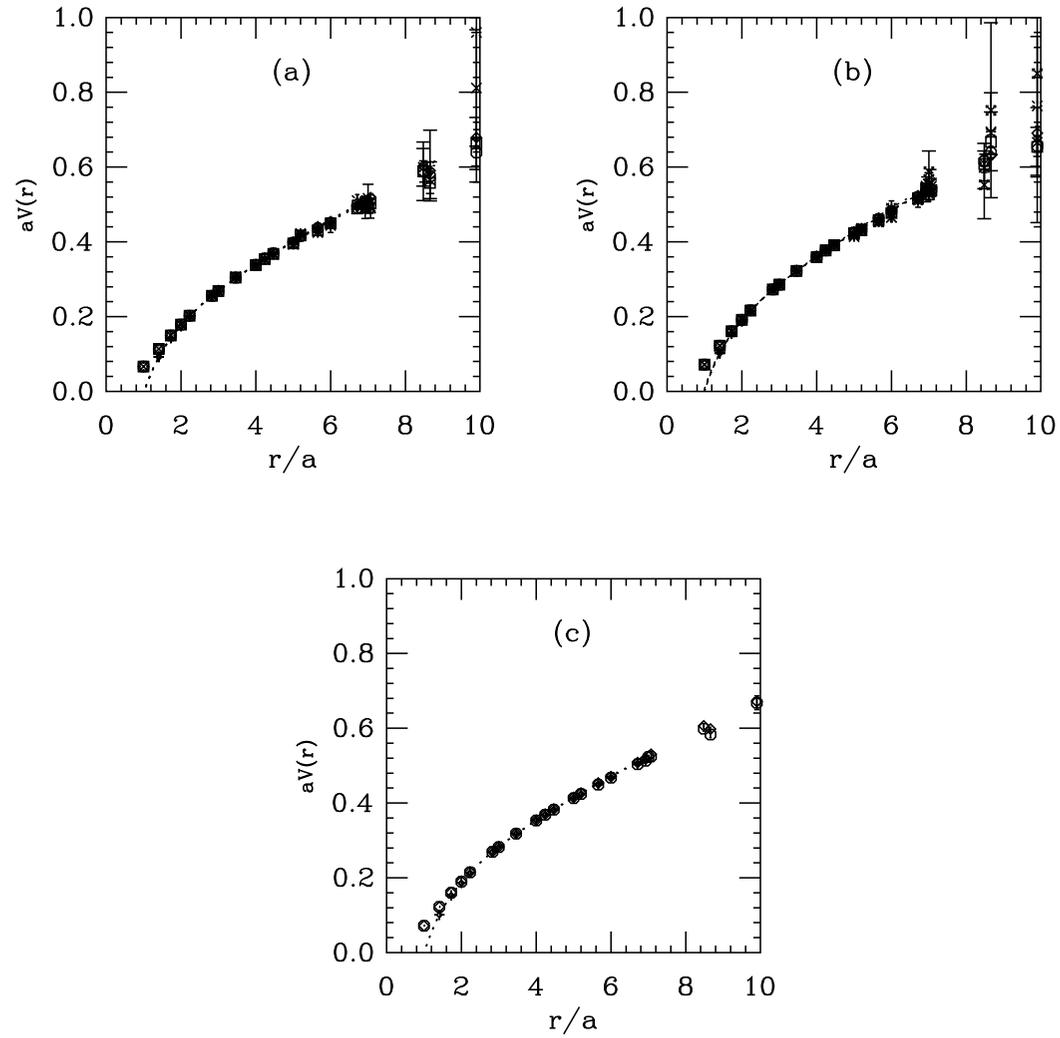
But – large N isn't about small m_q or even $a \rightarrow 0$, so it's easy

Potentials set the scale

- $16^3 \times 32$ volume
- Match Sommer parameter, $r_0 = 0.5$ fm ($r_0^2 F(r_0) = 1.65$) or $r_1 = 0.3$ fm
- Lattice spacing $a = 0.08$ fm or $1/a = 2400$ MeV
- $r_0\sqrt{\sigma} = 0.5$ fm \times 440 MeV = 1.12

Fubdamental rep fermion data sets

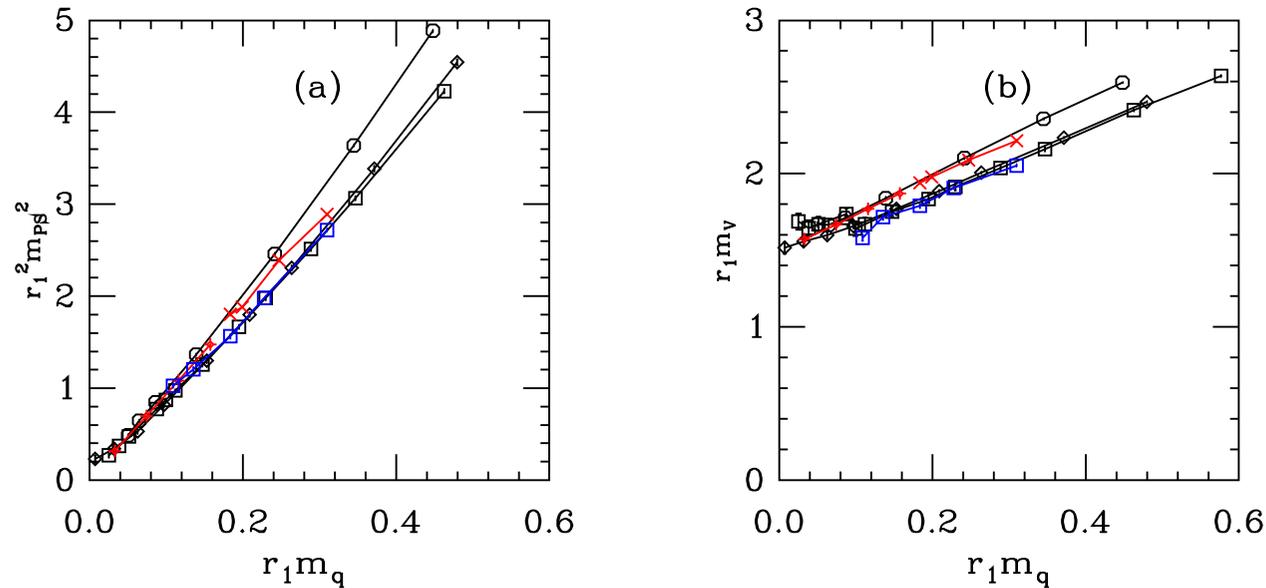
N	3	5	7
lattices	80	120	160
β	6.0175	17.5	34.9
$\lambda = g^2 N$	2.99	2.85	2.80
r_1/a	3.88(2)	3.77(2)	3.89(2)
$r_0\sqrt{\sigma}$	1.174(4)	1.172(3)	1.168(2)



Potentials for $SU(3)$, $SU(5)$, $SU(7)$. Can you tell which is which?

Meson spectroscopy

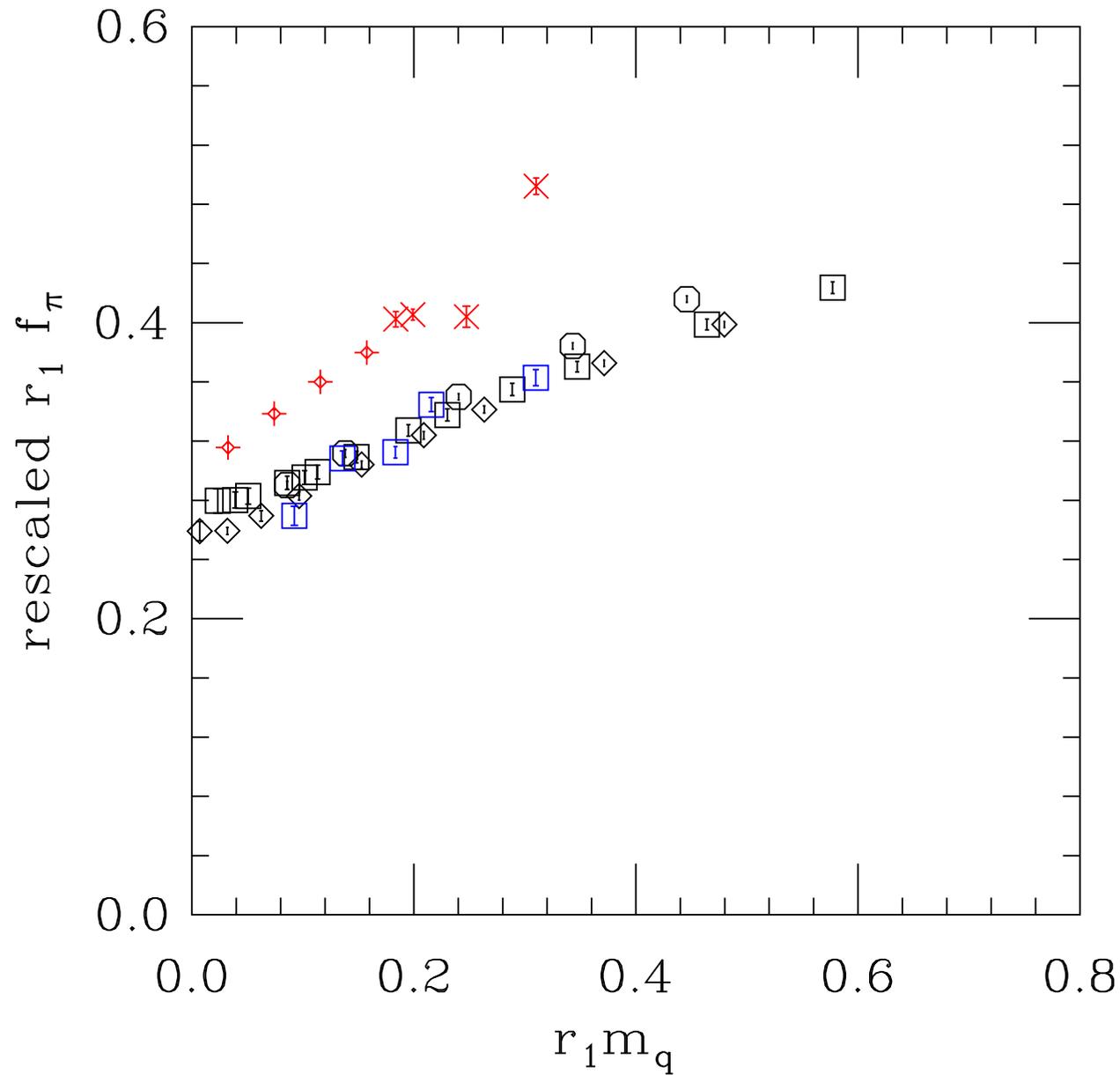
Meson masses should be N -independent



N -independence, also seen by everybody else, over much wider N range (Bali et al) (a) pseudoscalar mass squared (b) vector ($1/r_1 \sim 650$ MeV)

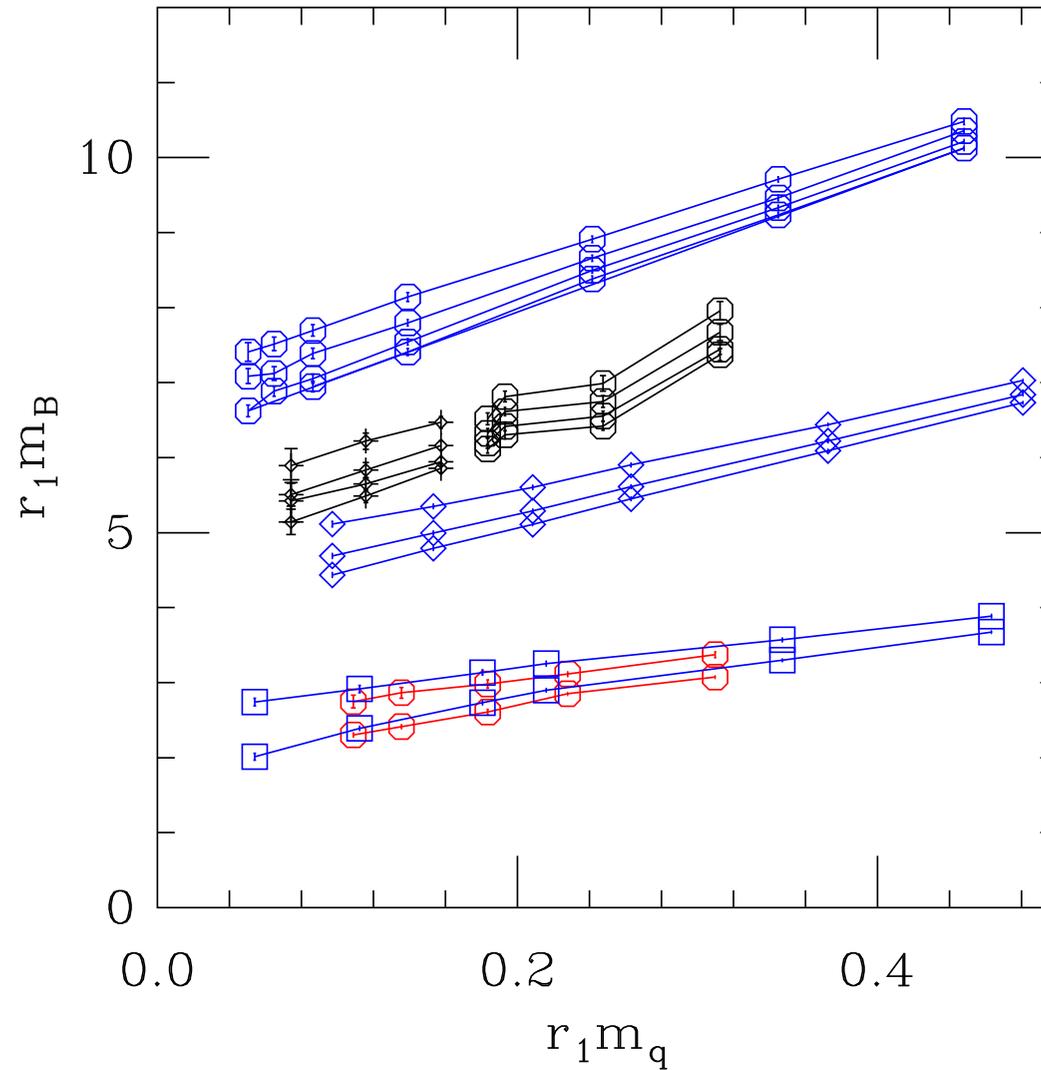
Black squares, diamonds, octagons for quenched $SU(3)$, $SU(5)$, $SU(7)$ fundamentals, red crosses for $SU(4)$ AS2; the fancy diamonds are the PQ data. Blue squares $SU(3)$ with two dynamical, fundamental flavors.

Decay constants $\langle 0|V|h\rangle \propto 1/\sqrt{N} \times N \propto \sqrt{N}$ for fundamentals, $\propto N$ for AS2



Pseudoscalar decay constant scaled to $SU(3)$ fundamentals by dividing by appropriate power of $N/3$

Baryon spectroscopy



The blue data are from the top quenched $SU(7), SU(5)$ and $SU(3)$ data. The red octagons are $SU(3)$ with dynamical fermions. The black lines are the six-quark baryons in $SU(4)$ AS2, octagons for dynamical and fancy diamonds for partially quenched.

Baryons – theory

The Rotor spectrum: large N baryon masses generically obey

$$M(N, J) = Nm_0 + \frac{J(J+1)}{N}B + \dots \quad (1)$$

m_0 and B are m_q dependent, so do comparisons versus m_q

Test the $J(J+1)$: ratios of mass differences are pure numbers

$$\Delta(J_1, J_2, J_3) = \frac{M(N, J_2) - M(N, J_3)}{M(N, J_1) - M(N, J_3)}, \quad (2)$$

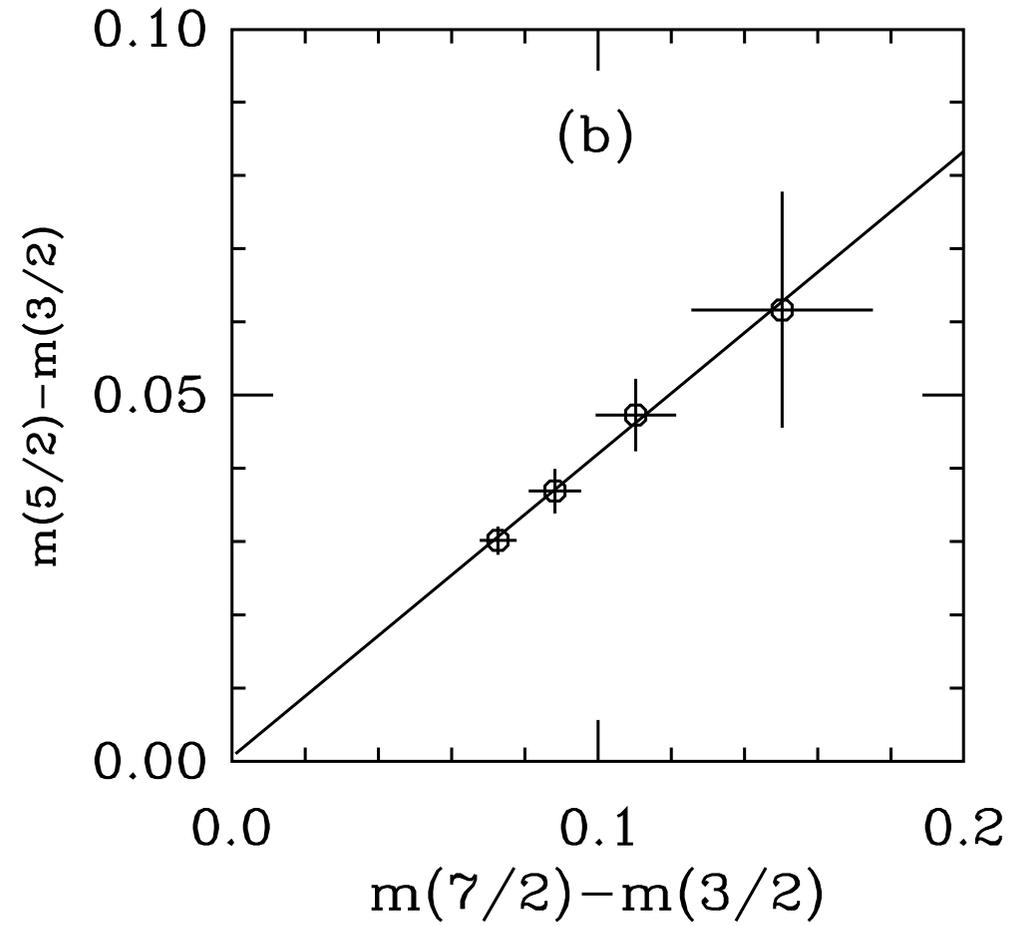
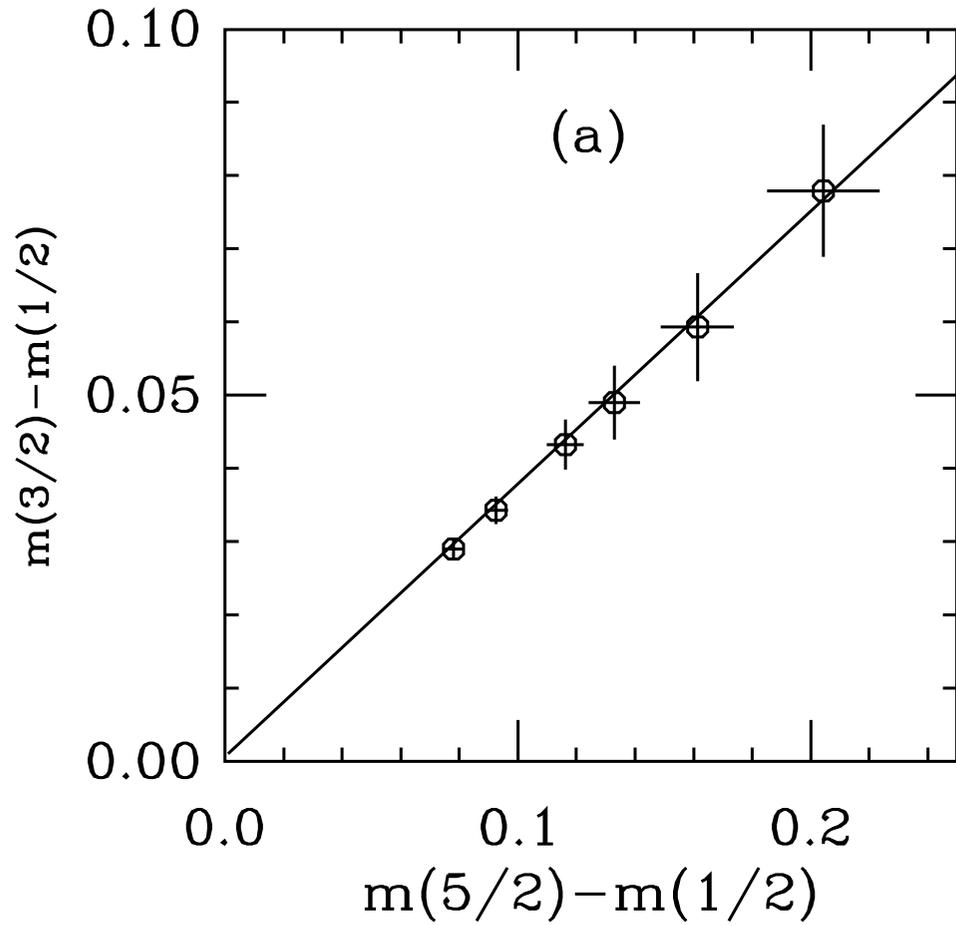
Exposing B : from the bottom of the multiplet

$$M(N, 3/2) - M(N, 1/2) = \frac{3B}{N} \quad (3)$$

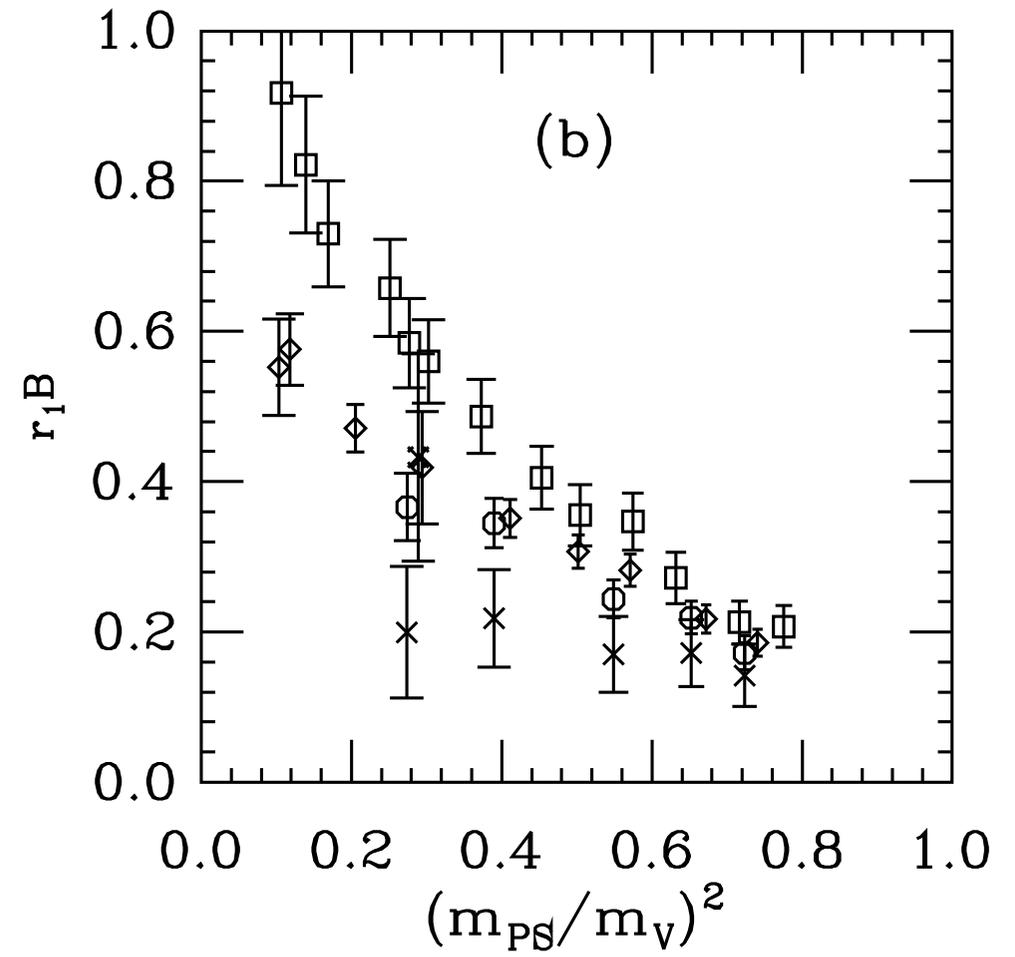
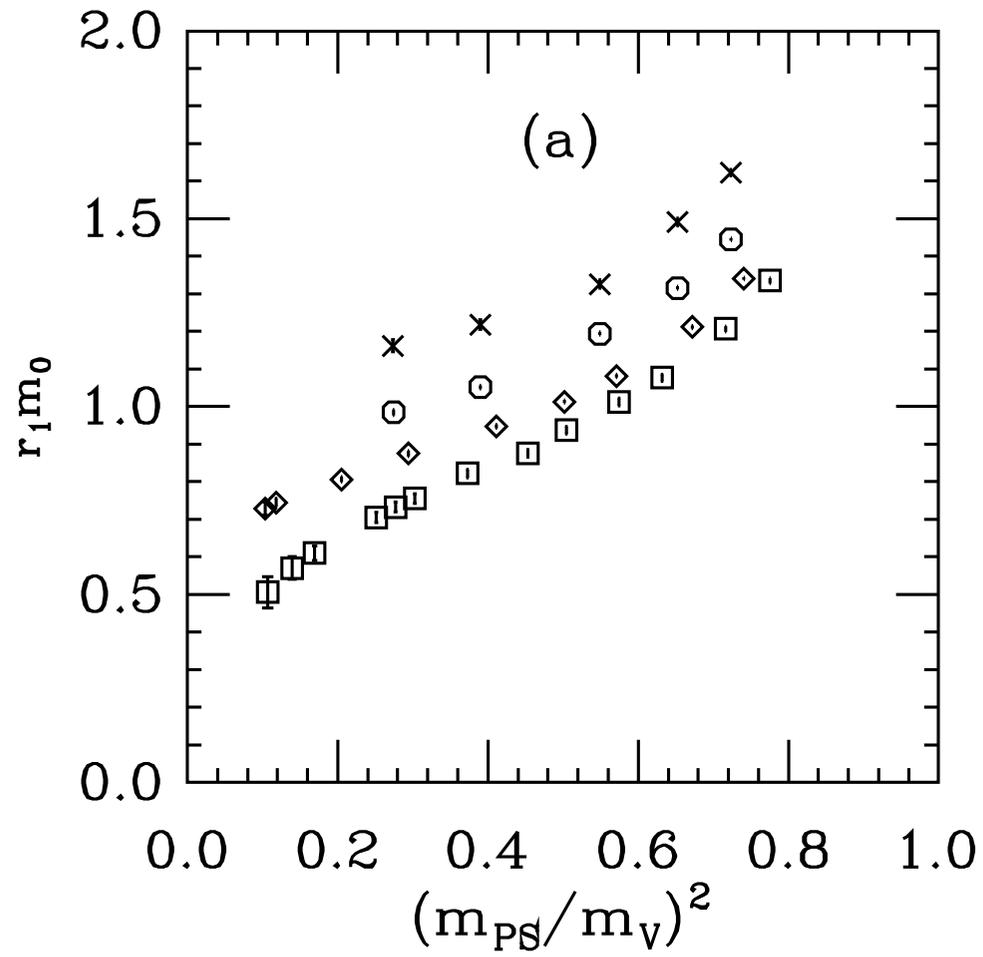
From the top of the multiplet a “rescaled Landé interval rule”

$$M(N, J = N/2) - M(N, J = N/2 - 1) = B \quad (4)$$

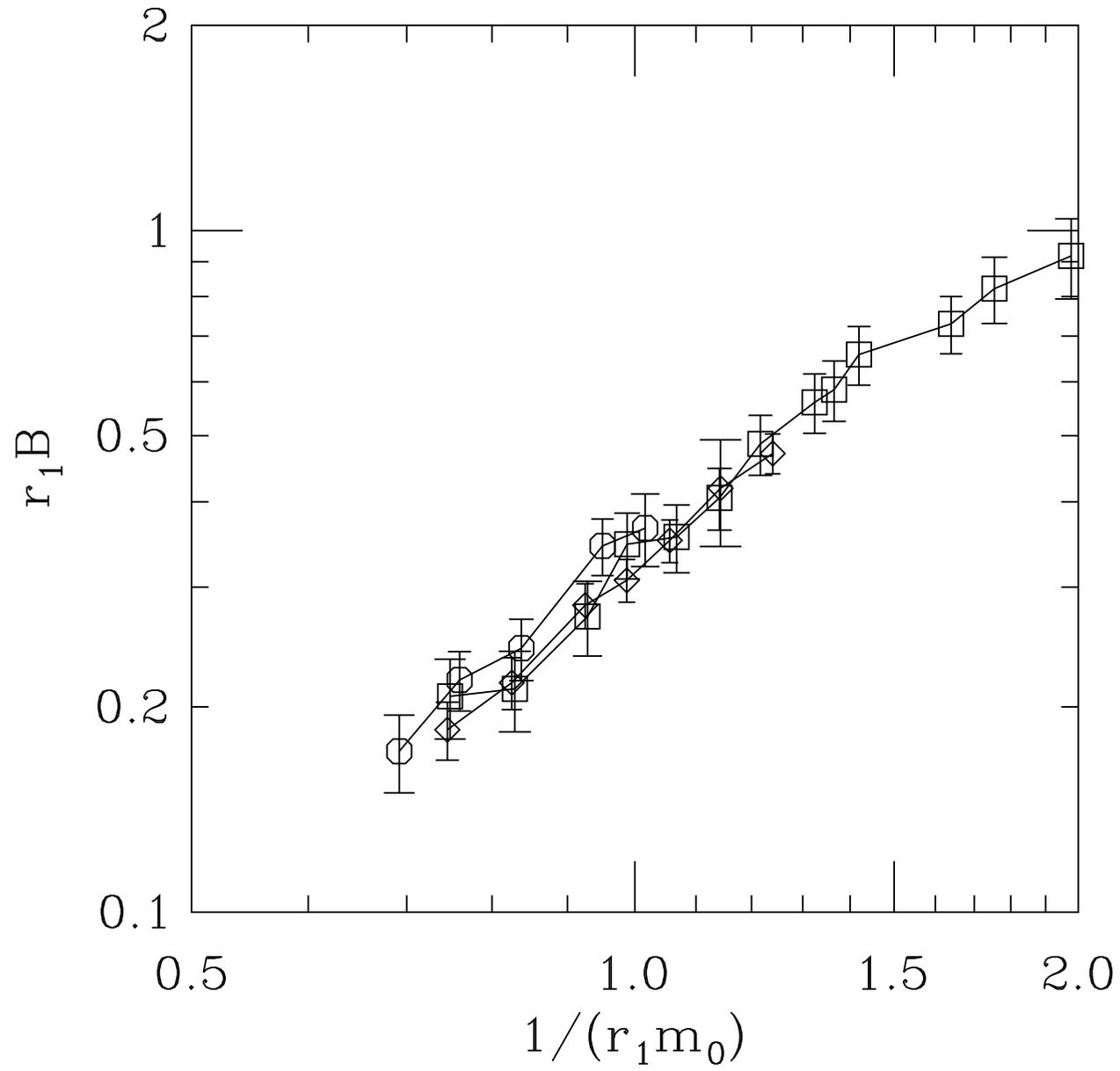
Or, just fit all states’ masses to two parameters (m_0 and B), plot vs m_q



Ratios of mass differences for $N = 5$ and $N = 7$. Lines are the analytic ratio (NOT a fit)



m_0 and B from $N = 3, 5, 7$. Crosses extrapolations to $1/N \rightarrow 0$. Modest drift with $1/N$.



B vs $1/m_0$ from quenched $SU(N)$ fundamental quark baryons

On to flavor $SU(3)$

The obvious project: flavor $SU(3)$ spectra, with different nonstrange and strange masses

- Lots of mass relations in $1/N$ and $\delta m = m_s - m_{ns}$
- Most published predictions are only for $J = 1/2$ and $3/2$ (because those are the only states we have)

There are lots of states, analogs of Δ 's, Σ 's, Λ 's, nucleons at many J 's

The project was mostly an issue of organization (and CG coefficients)

A quark model story for baryon spectroscopy

Colorspin! (De Rujula, George, Glashow; also MIT bag model, 1974)

m_i is the constituent quark mass

$$H_0 = \sum_i n_i m_i \quad (5)$$

There is also a color magnetic dipole (hyperfine) interaction

$$H_1 = \sum_{i \neq j} F_{ij} \lambda_i \lambda_j \sigma_i \sigma_j \quad (6)$$

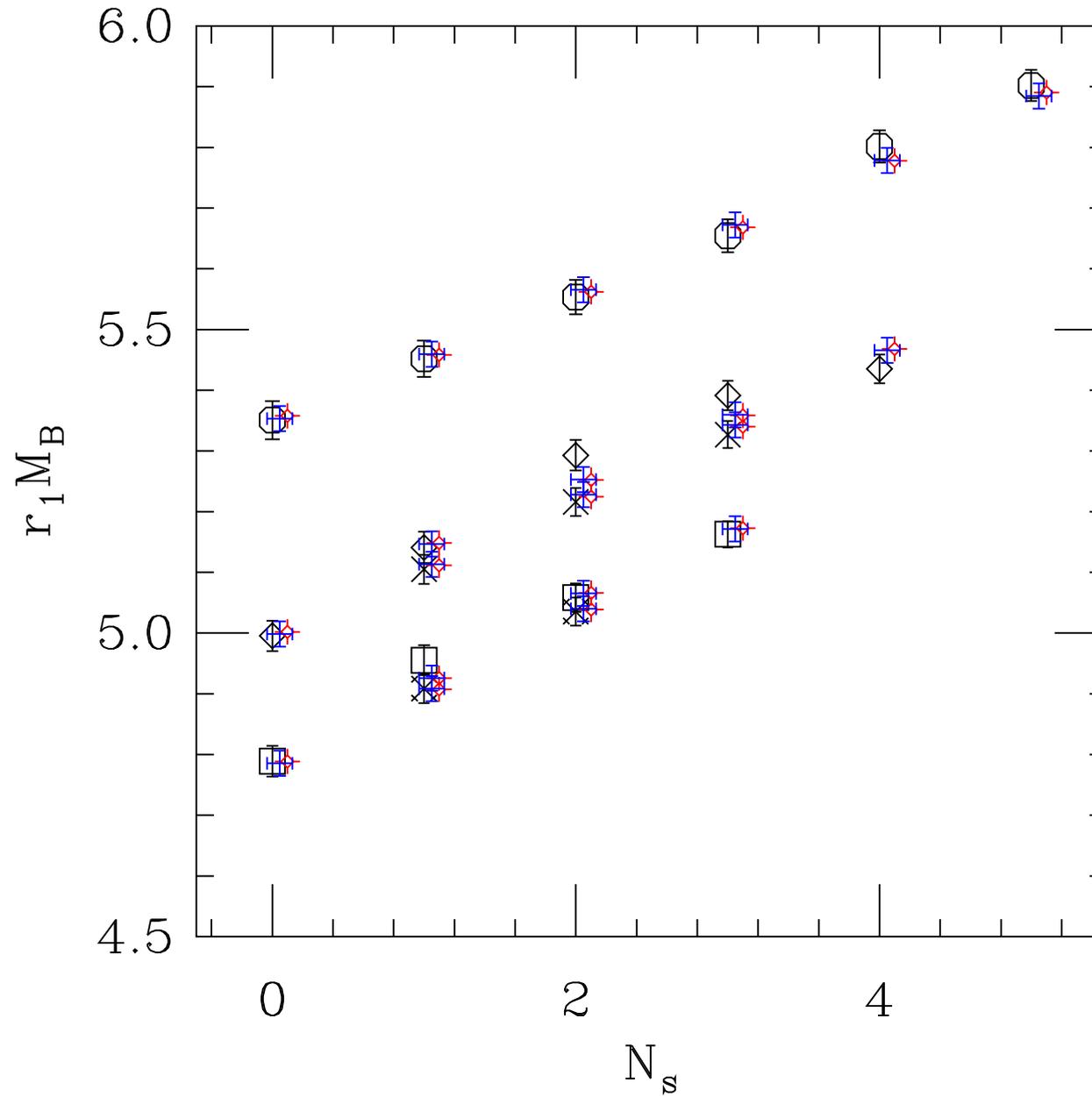
where

$$F_{ij} = \mu_i \mu_j = \frac{g}{m_i} \frac{g}{m_j} \propto g^2 = \frac{\lambda}{N} \quad (7)$$

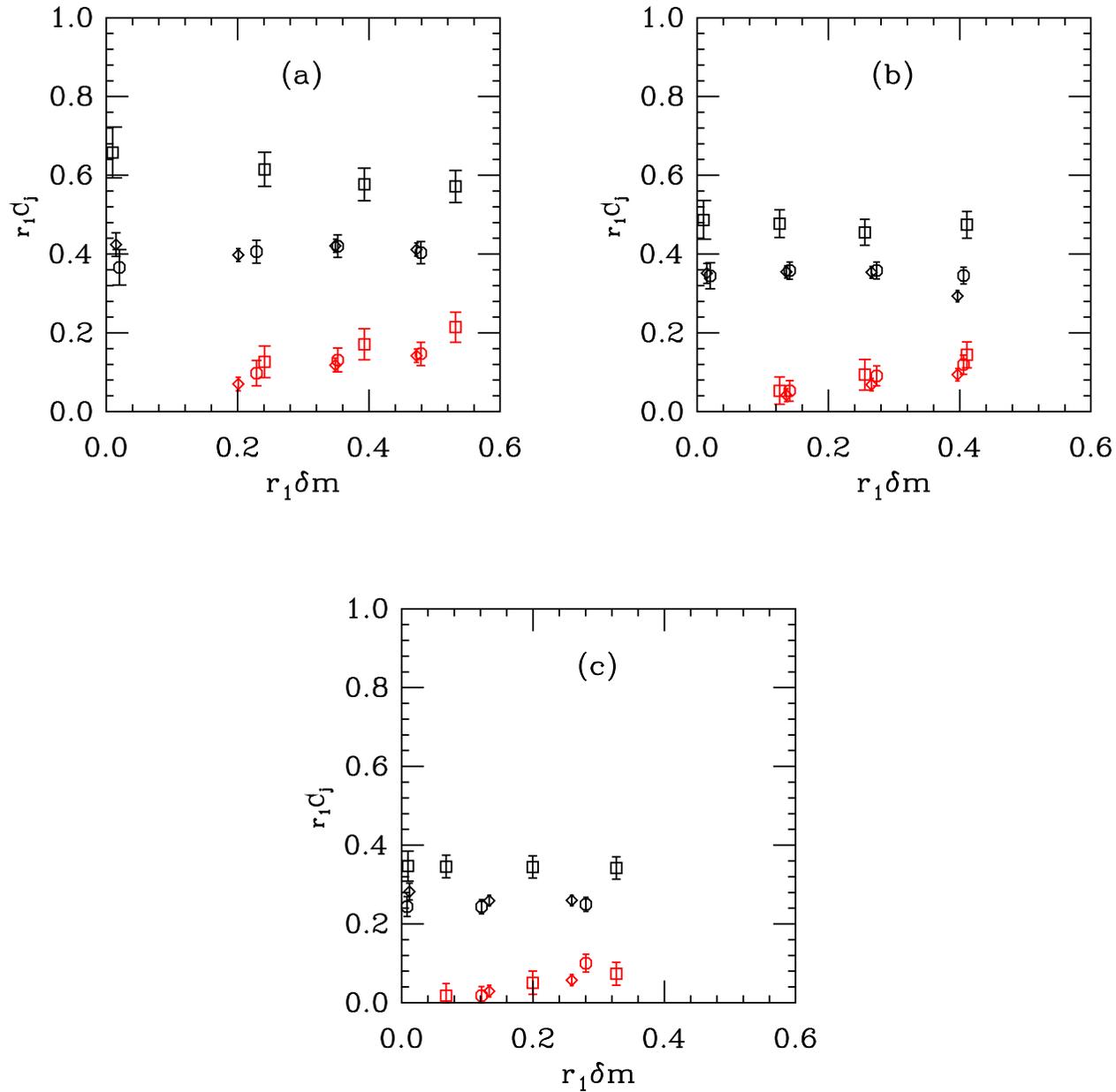
For degenerate mass quarks the rotor spectrum follows, simply,

$$\frac{f}{N} \left(\sum_i \sigma_i \sum_j \sigma_j - \sum_i \sigma_i^2 \right) = f \frac{J^2}{N} + C \quad (8)$$

For flavor $SU(3)$, slightly increase $m_s = m_i + \delta m$ and decrease the HFS term between s quarks



$SU(5)$ spectrum at one set of quark masses with two large- N fits overlaid



Color hyperfine interaction model parameters from fits to flavor $SU(3)$ data at matched $(m_{PS}/m_V)^2$. Data are squares for $N = 3$, diamonds for $N = 5$ and octagons for $N = 7$.

Summary

Surprisingly easy to see interesting physics

- Meson masses show expected large- N behavior
- Meson decay constants (mostly) scale appropriately with N
- Rotor spectrum for baryons works well
- Can see corrections to leading order, they seem to be $1/N$ with unsurprising size
- Flavor $SU(3)$ works as expected

Properties of mesons and baryons seem quite generic

Still things to do:

- Smaller m_q , to test chiral PT – this will require dynamical fermions
- Could be a large program of matrix elements vs $1/N$
- Large N will play a big role in the presentation of our (Neil, Svetitsky, Shamir, Jay, Liu) composite Higgs studies

It's a fun project, there is room for more people–

Thanks, Mike, for finding such a nice playground!